

**TIP O'ZGARISH CHIZIG'I SILLIQ BO'LMAGAN PARABOLIK-GIPERBOLIK
TENGLAMA UCHUN INTEGRAL ULASH SHARTLI CHEGARAVIY MASALA**

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Abstract: Tip o'zgarish chizig'i silliq bo'lmagan parabolik-giperbolik tenglama uchun integral ulash shartli chegaraviy masalaning bir qiymatli yechilishi isbotlangan. Masala yechimining yagonaligi energiya integrallari usulida, mavjudligi esa integral tenglamalar usulida isbotlangan.

Keywords: parabolik-giperbolik tenglama; integral ulash sharti; energiya integrallari; Grin funksiyasi; Dalamber formulasi; integral tenglama

About: FARS Publishers has been established with the aim of spreading quality scientific information to the research community throughout the universe. Open Access process eliminates the barriers associated with the older publication models, thus matching up with the rapidity of the twenty-first century.

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**BOUNDARY PROBLEM WITH INTEGRAL GLUING CONDITION FOR
PARABOLIC-HYPERBOLIC EQUATION WITH NON-SMOOTH LINE OF
TYPE-CHANGING**

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Abstract: A unique solvability of a boundary problem with integral gluing condition for parabolic-hyperbolic equation with non-smooth line of type-changing has been proved. The uniqueness of the solution to the problem is proved by energy integral's method, the existence by the method of integral equations..

Keywords: Parabolic-hyperbolic equation; integral gluing condition; energy integrals; Green's function; d'Alambert's formula; integral equations

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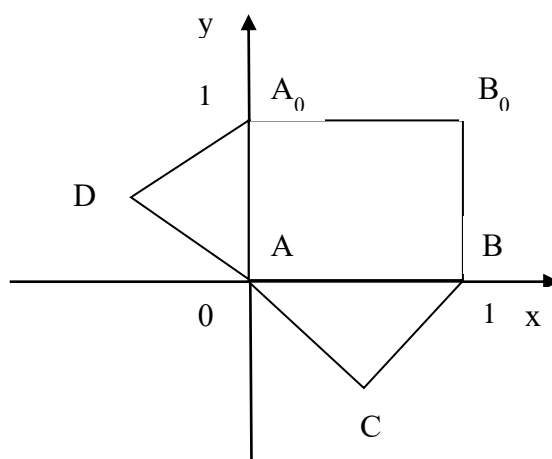
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Bizga Ω sohada quyidagi tenglama berilgan bo'lsin:

$$\begin{cases} U_{xx} - U_y = 0, & (x, y) \in \Omega_0 \\ U_{xx} - U_{yy} = 0, & (x, y) \in \Omega_i, (i = 1, 2) \end{cases} \quad (1)$$

Ω soha esa quyidagicha $\Omega = \Omega_0 \cup \Omega_1 \cup \Omega_2 \cup AB \cup AA_0$, bu yerda

$$A(0,0); A_0(0,1); B(1,0); B_0(1,1); C\left(\frac{1}{2}, -\frac{1}{2}\right); D\left(-\frac{1}{2}, \frac{1}{2}\right) \text{ to'g'ri}$$



1-rasm

Masala. (1) tenglamani qanoatlantiruvchi Ω sohada shunday $U(x, y)$ funksiya topilsinki, u $U(x, y) \in C_{x,y}^{2,1}(\Omega_0) \cap C(\bar{\Omega}) \cap C^2(\Omega_i), (i=1,2)$ regulyarlik shartlarini hamda quyidagi shartlarni qanoatlantirsin:

$$U(x, y)|_{AC} = \psi_1(x), \quad 0 \leq x \leq \frac{1}{2}, \quad (2)$$

$$U(x, y)|_{AD} = \psi_2(y), \quad -\frac{1}{2} \leq x \leq 0, \quad (3)$$

$$U(x, y)|_{BB_0} = \varphi(y), \quad 0 \leq y \leq 1, \quad (4)$$

$$U_y(x, +0) = I_1(U_y(x, -0)), \quad U_y(x, +0) = U_y(x, -0), \quad 0 < x < 1, \quad (5)$$

$$U_x(+0, y) = I_2(U_x(-0, y)), \quad U_x(+0, y) = U_x(-0, y), \quad 0 < y < 1. \quad (6)$$

Bu yerda I_1 va I_2 lar hozircha ixtiyoriy integral operatorlar.

Ω_0 sohada $U_{xx} - U_y = 0$ tenglama uchun birinchi chegaraviy masala yechimini quyidagicha yozamiz [1]:

$$u(x, y) = \int_0^1 \tau_1^+(\xi) G_1(x, y; \xi, 0) d\xi + \int_0^y \tau_2^+(\eta) G_{1\xi}(x, y; 0, \eta) d\eta - \int_0^y \varphi(\eta) G_{1\xi}(x, t; l, \eta) d\eta$$

bu yerda, $\tau_1^+(x) = U(x, 0), \quad 0 \leq x \leq 1,$

$$\tau_2^+(y) = U(0, y), \quad 0 \leq y \leq 1, \quad \varphi(y) = U(1, y), \quad 0 \leq y \leq 1.$$

Ω_1 sohada $U_{xx} - U_{yy} = 0$ tenglamaning $U(x, -0) = \tau_1^-(x)$, $0 \leq x \leq 1$,
 $U_y(x, -0) = \nu_1^-(x)$, $0 < x < 1$ boshlang'ich shartlarni qanoatlantiruvchi yechimini
Dalamber formulasi orqali yozamiz [2]:

$$U(x, y) = \frac{1}{2} [\tau_1^-(x-y) + \tau_1^-(x+y)] + \frac{1}{2} \int_{x-y}^{x+y} \nu_1^-(z) dz.$$

Bu yechimni (2) shartga bo'ysundiramiz:

$$2\psi_1(x) = \tau_1^-(0) + \tau_1^-(2x) + \int_{2x}^0 \nu_1^-(z) dz, \quad 0 \leq x \leq \frac{1}{2}.$$

Endi $2x$ ni x bilan almashtirsak

$$2\psi_1\left(\frac{x}{2}\right) = \tau_1^-(0) + \tau_1^-(x) - \int_0^x \nu_1^-(z) dz, \quad 0 \leq x \leq 1$$

hosil bo'ladi. Keyin bu tenglikni x bo'yicha differensiallasak

$$\psi_1'\left(\frac{x}{2}\right) = \tau_1'(x) - \nu_1^-(x), \quad 0 < x < 1 \quad (7)$$

Munosabatni olamiz.

Ω_2 sohada $U_{xx} - U_{yy} = 0$ tenglamaning $U(-0, y) = \tau_2^-(y)$, $0 \leq y \leq 1$,
 $U_x(-0, y) = \nu_2^-(y)$, $0 < y < 1$ boshlang'ich shartlarni qanoatlantiruvchi yechimini
Dalamber formulasi orqali yozamiz [2]:

$$U(x, y) = \frac{1}{2} [\tau_2^-(y-x) + \tau_2^-(y+x)] - \frac{1}{2} \int_{y+x}^{y-x} \nu_2^-(z) dz.$$

Bu yechimni (3) shartga bo'ysundiramiz:

$$2\psi_2(y) = \tau_2^-(2y) + \tau_2^-(0) - \int_0^{2y} \nu_2^-(z) dz, \quad 0 \leq y \leq \frac{1}{2}.$$

Endi $2y$ ni y bilan almashtirsak,

$$2\psi_2\left(\frac{y}{2}\right) = \tau_2^-(y) + \tau_2^-(0) - \int_0^y \nu_2^-(z) dz, \quad 0 \leq y \leq 1$$

hosil bo'ladi. Keyin bu tenglikni y bo'yicha differensiallasak

$$\psi_2'\left(\frac{y}{2}\right) = \tau_2'(y) - \nu_2^-(y), \quad 0 < y < 1 \quad (8)$$

tenglikka ega bo'lamiz.

Masala yechimining yagonaligi.

Masala yechimi 2 ta u_1 va u_2 funksiyalar bo'lsin. U holda $v = u_1 - u_2$ funksiyaga nisbatan quyidagi masalani qarashimiz mumkin bo'ladi:

$$\begin{cases} v_{xx} - v_y = 0, & (x, y) \in \Omega_0 \\ v_{xx} - v_{yy} = 0, & (x, y) \in \Omega_i, (i=1,2) \end{cases}$$

$$v(x, y)|_{AC} = 0, v(x, y)|_{AD} = 0, v(x, y)|_{BB_0} = 0,$$

$$v_y(x, +0) = I_1(v_y(x, -0)), v_x(+0, y) = I_2(v_x(-0, y))$$

Bu esa bir jinsli tenglama uchun bir jinsli shartli masaladir. Agar ushbu masala yechimining $v(x, y) \equiv 0$ ekanligini ko'rsata olsak, asosiy masalaning yechimi yagonaligi kelib chiqadi:

Dastlab, $v_{xx} - v_y = 0$ tenglamada $y \rightarrow +0$ holda limitga o'tamiz va

$$\tau_1^{+''}(x) - v_1^+(x) = 0$$

tenglamani olamiz. Bu yerda $\tau_1^+(x) = v(x, +0)$, $v_1^+(x) = v_y(x, +0)$ hamda ushbu tenglamani $\tau_1^+(x)$ ga ko'paytirib (0,1) oraliqda integrallaymiz:

$$\int_0^1 \tau_1^+(x) \tau_1^{+''}(x) dx - \int_0^1 \tau_1^+(x) v_1^+(x) dx = 0.$$

Ushbu integralni bo'laklab integrallaymiz:

$$\tau_1^+(x) \tau_1^{+'}(x) \Big|_0^1 - \int_0^1 \left[\tau_1^{+'}(x) \right]^2 dx - \int_0^1 \tau_1^+(x) v_1^+(x) dx = 0.$$

$\tau_1^+(0) = \tau_1^+(1) = 0$ ekanligini inobatga olsak yuqoridagi ifoda quyidagi ko'rinishga keladi:

$$\int_0^1 \left[\tau_1^{+'}(x) \right]^2 dx + \int_0^1 \tau_1^+(x) v_1^+(x) dx = 0$$

Bundan $\int_0^1 \left[\tau_1^{+'}(x) \right]^2 dx \geq 0$ bo'lgani uchun $\int_0^1 \tau_1^+(x) v_1^+(x) dx \geq 0$ ekanligini ko'rsata olsak yechim yagonaligini ko'rsatishimiz mumkin bo'ladi. Buning uchun I_1 operatorni quyidagicha tanlaymiz:

$$I_1(f(x)) = \int_0^x f(z) K(x, z) dz.$$

Demak,

$$\begin{aligned}
 J_1 &= \int_0^1 \tau_1^+(x) \nu_1^+(x) dx = \int_0^1 \tau_1^+(x) I_1(\tau_1^+(x)) dx = \int_0^1 \tau_1^+(x) dx \int_0^x \tau_1^+(z) K(x, z) dz = \\
 &= \int_0^1 \tau_1^+(x) dx \left[\tau_1^+(z) K(x, z) \Big|_0^x - \int_0^x \tau_1^+(z) \frac{\partial}{\partial z} K(x, z) dz \right] = \\
 &= \int_0^1 \tau_1^2(x) K(x, x) dx - \int_0^1 \tau_1(x) dx \int_0^x \tau_1^+(z) \frac{\partial}{\partial z} K(x, z) dz.
 \end{aligned}$$

Bundan $K(x, x) \geq 0$ deb olib, $K(x, z)$ uchun quyidagi tenglikni o'rinli deb olsak

$$\frac{\partial}{\partial z} K(x, z) = -K_1(x) K_1(z),$$

J_1 integralimizni quyidagicha ifodalashimiz mumkin bo'ladi:

$$\begin{aligned}
 &\int_0^1 \tau_1^2(x) K(x, x) dx + \frac{1}{2} \int_0^1 \frac{d}{dx} \left[\int_0^x \tau_1^+(z) K_1(z) dz \right]^2 dx = \\
 &= \int_0^1 \tau_1^2(x) K(x, x) dx + \frac{1}{2} \left(\int_0^1 \tau_1^+(z) K_1(z) dz \right)^2.
 \end{aligned}$$

Shunday qilib yuqoridagi shartlarni qanoatlantiruvchi $K(x, z)$ funksiya uchun $\tau_1^-(x)$ funksiyani nolga tengligi kelib chiqadi. Bundan esa $\nu_1^-(x)$ funksiyani ham nolga tengligi kelib chiqadi.

Agar $\tau_1^-(x) = 0, \nu_1^-(x) = 0$ bo'lsa D'alamber formulasidan $v(x, y)$ ni $v(x, y) \equiv 0$ ekanligi kelib chiqadi, bu esa bizga yechim yagonaligini keltirib chiqaradi.

Endi $v_{xx} - v_y = 0$ tenglamani olib uni v ga ko'paytirib Ω_0 soha bo'yicha integrallaymiz:

$$\begin{aligned}
 &\iint_{\Omega_0} v(v_{xx} - v_y) dx dy = 0 \text{ yoki} \\
 &\iint_{\Omega_0} \left[(vv_x)_x - v_x^2 - \frac{1}{2}(v^2)_y \right] dx dy = 0.
 \end{aligned}$$

Grin formulasini qo'llab [2], soha bo'yicha integralni soha chegarasi bo'yicha integralga keltiramiz:

$$\int_{\partial\Omega_0} vv_x dy + \frac{1}{2}(v^2) dx - \iint_{\Omega_0} v_x^2 dx dy = 0.$$

Endi chegara bo'yicha integrallashni amalga oshirib olamiz:

$$\int_0^1 \frac{1}{2} [\tau_1^+(x)]^2 dx + \int_0^1 v(1, y) v_x(1, y) dy - \frac{1}{2} \int_0^1 v^2(x, 1) dx - \int_0^1 \tau_2^+(y) v_2^+(y) dy - \iint_{\Omega_0} v_x^2(x, y) dx dy = 0.$$

$\tau_1^+(x) = 0, v(1, y) = 0$ ekanligidan integral quyidagi ko'rinishga keladi:

$$\frac{1}{2} \int_0^1 v^2(x, 1) dx + \int_0^1 \tau_2^+(y) v_2^+(y) dy + \iint_{\Omega_0} v_x^2(x, y) dx dy = 0.$$

Agar $\int_0^1 \tau_2^+(y) v_2^+(y) dy \geq 0$ ni ekanligini ko'rsata olsak yechim yagonaligini

ko'rsatishimiz mumkin.

$$\begin{aligned} J_2 &= \int_0^1 \tau_2^+(y) v_2^+(y) dy = \int_0^1 \tau_2^+(y) I_2(\tau_2'(y)) dy = \left\{ I_2(f(x)) = \int_0^x f(x) P(x, z) dz \right\} = \\ &= \int_0^1 \tau_2^+(y) \left[\int_0^y \tau_2'(z) P(y, z) dz \right] dy = \int_0^1 \tau_2^+(y) \left[\tau_2^+(z) P(y, z) \Big|_0^y - \int_0^y \tau_2(z) \frac{\partial}{\partial z} P(y, z) dz \right] dy = \\ &= \int_0^1 (\tau_2^+(y))^2 P(y, y) dy - \int_0^1 \tau_2^+(y) \left[\int_0^y \tau_2(z) \frac{\partial}{\partial z} P(y, z) dz \right] dy = \\ &= \left\{ \frac{\partial}{\partial z} P(y, z) = -P_1(y) P_1(z) \right\} = \int_0^1 (\tau_2^+(y))^2 P(y, y) dy + \\ &+ \int_0^1 \tau_2^+(y) \left[\int_0^y \tau_2(z) P_1(y) P_1(z) dz \right] dy = \int_0^1 (\tau_2^+(y))^2 P(y, y) dy + \\ &+ \frac{1}{2} \int_0^1 \frac{d}{dy} \left[\int_0^y \tau_2^+(z) P_1(z) dz \right]^2 dy = \int_0^1 (\tau_2^+(y))^2 P(y, y) dy + \frac{1}{2} \left[\int_0^1 \tau_2^+(z) P_1(z) dz \right]^2 \geq 0. \end{aligned}$$

Agar $P(y, y) \geq 0$ va $\frac{\partial}{\partial z} P(y, z) = -P_1(y) P_1(z)$ larni ta'minlay olsak $P(y, z)$

funksiya uchun $\tau_2(y)$ funksiyani nolga tengligi kelib chiqadi. Bundan esa $v_2(y)$ funksiyani ham nolga tengligi kelib chiqadi.

Agar $\tau_2(y) = 0, v_2(y) = 0$ bo'lsa D'alamber formulasida $v(x, y)$ ni $v(x, y) \equiv 0$ ekanligi kelib chiqadi, bu esa bizga yechim yagonaligini keltirib chiqaradi.

Masala yechimining mavjudligi.

Dastlab (5) va (6) shartlarni quyidagicha ko'rishga keltirib olamiz:

$$v_1^+(x) = I_1(v_1^-(x)), \quad (5)$$

$$v_2^+(x) = I_2(v_2^-(x)). \quad (6)$$

Ω_0 sohada $y \rightarrow +0$ da limitga o'tgan ifodamiz

$$\tau_1^{+''}(x) - v_1^+(x) = 0 \quad (9)$$

va Ω_1 sohada olgan

$$\psi_1'\left(\frac{x}{2}\right) = \tau_1^-(x) - v_1^-(x), \quad 0 < x < 1 \quad (7)$$

munosabatlarimiz hamda $\tau_1^-(0) = \psi_1(0)$, $\tau_1^-(1) = \varphi(0)$ shartlarga asosan quyidagi masalaga kelamiz:

$$\begin{cases} \tau_1^{+''}(x) - I_1(\tau_1^+(x)) = -I_1\left(\psi_1'\left(\frac{x}{2}\right)\right) \\ \tau_1^+(0) = \psi_1(0), \quad \tau_1^+(1) = \varphi(0). \end{cases}$$

$K(x, z)$ yadroni $K(x, z) = K_1(x)K_1(z)$ deb olsak I_1 operatorni

$$I_1(f(x)) = \int_0^x f(z)K_1(x)K_1(z)dz \quad (10)$$

ko'rishga keladi. Natijada quyidagi integral tenglamaga kelamiz:

$$\tau_1^{+''}(x) - \int_0^x \tau_1^+(z)K_1(x)K_1(z)dz = \psi_1(x), \quad (11)$$

bu yerda $\psi_1(x) = -\int_0^x \psi_1'\left(\frac{x}{2}\right)K_1(x)K_1(z)dz$.

Hosil bo'lgan (11) integral tenglamani $(0, x)$ da integrallab

$$\int_0^x \tau_1^{+''}(\xi)d\xi - \int_0^x K_1(\xi)d\xi \int_0^\xi \tau_1^+(z)K_1(z)dz = \int_0^x \psi_1(\xi)d\xi$$

ifodaga ega bo'lamiz va

$$\int_0^x \tau_1^+(\xi)d\xi = \tau_1^+(x) - \tau_1^+(0)$$

ni hisobga olib integrallash tartibini o'zgartirsak quyidagi integral tenglamaga kelamiz:

$$\tau_1^+(x)dx - \int_0^x \tau_1^+(z)K_1(z)dz \int_z^x K_1(\xi)d\xi = \psi_1(x) + \tau_1^+(0)$$

$$\text{bu yerda } \psi_1(x) = \int_0^x \psi_1(\xi) d\xi.$$

Hozircha $\tau_1'(0)$ ni ma'lum deb olib integral tenglamani yechishda davom etamiz va quyidagi belgilashlarni kiritib olamiz:

$$K_1(x, z) = K_1(z) \int_z^x K_1(\xi) d\xi - \text{yadro uzluksiz yoki kichik maxsuslikka ega bo'lishi}$$

mumkin.

$$f_1(x) = \psi_1(x) + \tau_1'(0) \text{ esa uzluksiz funksiya bo'lsin.}$$

Bularni hisobga olib integral tenglama yechimini rezolventa orqali quyidagicha yozish mumkin bo'ladi:

$$\tau_1'(x) = f_1(x) + \int_0^x R(x, z) f_1(z) dz.$$

Bu yerda $R(x, z)$ funksiya $K_1(x, z)$ yadroning rezolventasi.

Ushbu tenglikni $(0, x)$ intervalda integrallab $\tau_1(x)$ funksiyani topamiz.

$$\tau_1(x) - \tau_1(0) = \int_0^x f_1(\xi) d\xi + \int_0^x \int_0^\xi R(x, z) f_1(z) dz d\xi$$

yoki

$$\tau_1(x) = \int_0^x f_1(\xi) d\xi + \int_0^x \int_0^\xi R(x, z) f_1(z) dz d\xi + \psi_1(0).$$

Endi yuqorida ma'lum deb olib davom etgan qiymat $\tau_1'(0)$ ni topish uchun quyidagi hisob kitoblarni amalga oshiramiz:

$$\tau_1(x) = \int_0^x [\psi_1(\xi) + \tau_1'(0)] d\xi + \int_0^x \int_0^\xi R(x, z) [\psi_1(z) + \tau_1'(0)] dz d\xi + \psi_1(0) \quad (12)$$

yoki

$$\tau_1(x) = \int_0^x \psi_1(\xi) d\xi + \int_0^x \tau_1'(0) d\xi + \int_0^x \int_0^\xi R(x, z) \psi_1(z) dz d\xi + \int_0^x \int_0^\xi R(x, z) \tau_1'(0) dz d\xi + \psi_1(0).$$

$$\text{Ushbu belgilashni } F_1(x, z) = \int_0^x \psi_1(\xi) d\xi + \int_0^x \int_0^\xi R(x, z) \psi_1(z) dz d\xi + \psi_1(0) \text{ kiritib olib}$$

hisob kitobni sodda ko'rinishda tasvirlashimiz mumkin.

$$\tau_1(x) = \int_0^x \tau_1'(0) d\xi + \int_0^x \int_0^\xi R(x, z) \tau_1'(0) dz d\xi + F_1(x, z).$$

$x=1$ dagi $\tau_1(x)$ funksiyani qiymatidan foydalangan holda $\tau_1'(0)$ ni topishimiz mumkin bo'ladi:

$$\tau_1(1) = \tau_1'(0) + \tau_1'(0) \int_0^1 \int_0^\xi R(1, z) dz d\xi + \varphi(1, z),$$

$$\tau_1'(0) = \frac{\varphi(0) - \varphi(1, z)}{1 + \int_0^1 \int_0^\xi R(1, z) dz d\xi}.$$

Bundan $\int_0^1 \int_0^\xi R(1, z) dz d\xi \neq -1$ qo'shimcha shartga ega bo'lamiz.

$\tau_1(x)$ - (12) ko'rinishda topilgandan so'ng (7) tenglikdan foydalanib $v_1^+(x)$ ni topish mumkin bo'ladi.

Ulash shartiga ko'ra esa $v_1^-(x)$ aniq ko'rinishda topiladi.

Ω_2 sohada $U_{xx} - U_{yy} = 0$ tenglamaning $U(-0, y) = \tau_2^-(y)$, $0 \leq y \leq 1$, $U_x(-0, y) = v_2^-(y)$, $0 < y < 1$ boshlang'ich shartlarni qanoatlantiruvchi yechimini Dalamber formulasi orqali yozib olamiz [3].

$$U(x, y) = \frac{1}{2} [\tau_2(y-x) + \tau_2(y+x)] - \frac{1}{2} \int_{y+x}^{y-x} v_2^-(z) dz.$$

Bu yechimni (3) shartga bo'ysundirib olgan (8) tengligimizga I_2 operatorni qo'llab

$$I_2 \left(\tau_2'(y) \right) - I_2 \left[\psi_1' \left(\frac{y}{2} \right) \right] = v_2^+(y) \text{ tenglikka ega bo'lamiz.}$$

Keyin esa

$$\begin{aligned} v_2^+(y) = & \int_0^1 \tau_1(\xi) G_{1x}(0, y; \xi, 0) d\xi + \tau_2(\eta) \cdot \sum_{n=-\infty}^{+\infty} \frac{1}{\sqrt{y-\eta}} e^{-\frac{(2n)}{\sqrt{y-\eta}}} \Big|_0^y - \\ & - \int_0^y \tau_2'(\eta) \left(\sum_{n=-\infty}^{+\infty} \frac{1}{\sqrt{y-\eta}} e^{-\frac{(2n)}{\sqrt{y-\eta}}} \right) d\eta - \varphi(\eta) \cdot \sum_{n=-\infty}^{+\infty} \frac{1}{\sqrt{y-\eta}} e^{-\frac{(2n+1)}{\sqrt{y-\eta}}} \Big|_0^y + \\ & + \int_0^y \varphi'(\eta) \left(\sum_{n=-\infty}^{+\infty} \frac{1}{\sqrt{y-\eta}} e^{-\frac{(2n+1)}{\sqrt{y-\eta}}} \right) d\eta \end{aligned}$$

yoki

$$\begin{aligned}
 I_2\left(\tau_2'(y)\right) - I_2\left[\psi_1'\left(\frac{y}{2}\right)\right] &= \int_0^1 \tau_1(\xi)G_{1x}(0, y; \xi, 0)d\xi + \tau_2(\eta) \cdot \sum_{n=-\infty}^{+\infty} \frac{1}{\sqrt{y-\eta}} e^{-\frac{(2n)}{\sqrt{y-\eta}}} \Big|_0^y - \\
 - \int_0^y \tau_2'(\eta) \left(\sum_{n=-\infty}^{+\infty} \frac{1}{\sqrt{y-\eta}} e^{-\frac{(2n)}{\sqrt{y-\eta}}} \right) d\eta - \varphi(\eta) \cdot \sum_{n=-\infty}^{+\infty} \frac{1}{\sqrt{y-\eta}} e^{-\frac{(2n+1)}{\sqrt{y-\eta}}} \Big|_0^y + \\
 + \int_0^y \varphi'(\eta) \left(\sum_{n=-\infty}^{+\infty} \frac{1}{\sqrt{y-\eta}} e^{-\frac{(2n+1)}{\sqrt{y-\eta}}} \right) d\eta.
 \end{aligned}$$

Shundan so'ng quyidagiga ega bo'lamiz:

$$\begin{aligned}
 I_2\left(\tau_2'(y)\right) - I_2\left[\psi_1'\left(\frac{y}{2}\right)\right] &= \int_0^1 \tau_1(\xi)G_{1x}(0, y; \xi, 0)d\xi - \tau_2(0) \cdot \sum_{n=-\infty}^{+\infty} \frac{1}{\sqrt{y}} e^{-\frac{(2n)}{\sqrt{y}}} - \\
 - \int_0^y \tau_2'(\eta) \left(\sum_{n=-\infty}^{+\infty} \frac{1}{\sqrt{y-\eta}} e^{-\frac{(2n)}{\sqrt{y-\eta}}} \right) d\eta - \varphi(0) \cdot \sum_{n=-\infty}^{+\infty} \frac{1}{\sqrt{y}} e^{-\frac{(2n+1)}{\sqrt{y}}} + \\
 + \int_0^y \varphi'(\eta) \left(\sum_{n=-\infty}^{+\infty} \frac{1}{\sqrt{y-\eta}} e^{-\frac{(2n+1)}{\sqrt{y-\eta}}} \right) d\eta.
 \end{aligned}$$

Bu tenglikda ma'lum narsalarni bir tomonga olib o'tsak

$$I_2\left(\tau_2'(y)\right) + \int_0^y \tau_2'(\eta) \left(\sum_{n=-\infty}^{+\infty} \frac{1}{\sqrt{y-\eta}} e^{-\frac{(2n)}{\sqrt{y-\eta}}} \right) d\eta = \mathcal{P}_2^0(y) \quad (13)$$

hosil

bo'ladi.

Bu

yerda

$$\begin{aligned}
 \mathcal{P}_2^0(y) &= I_2\left[\psi_1'\left(\frac{y}{2}\right)\right] + \int_0^1 \tau_1(\xi)G_{1x}(0, y; \xi, 0)d\xi - \tau_2(0) \cdot \sum_{n=-\infty}^{+\infty} \frac{1}{\sqrt{y}} e^{-\frac{(2n)}{\sqrt{y}}} + \\
 + \varphi(0) \cdot \sum_{n=-\infty}^{+\infty} \frac{1}{\sqrt{y}} e^{-\frac{(2n+1)}{\sqrt{y}}} + \int_0^y \varphi'(\eta) \left(\sum_{n=-\infty}^{+\infty} \frac{1}{\sqrt{y-\eta}} e^{-\frac{(2n+1)}{\sqrt{y-\eta}}} \right) d\eta
 \end{aligned}$$

$$I_2 \text{ operatorni } I_2(g(y)) = \int_0^y g(y)P(y, z)dz \text{ ko'rinishda tanlab olib, } \quad (13)$$

tenglikni quyidagicha yozib olish mumkin:

$$\int_0^y \tau_2'(\eta) \left(P(y, \eta) + \sum_{n=-\infty}^{+\infty} \frac{1}{\sqrt{y-\eta}} e^{-\frac{(2n)}{\sqrt{y-\eta}}} \right) d\eta = \mathcal{P}_2^0(y). \quad (14)$$

(14) tenglikni y bo'yicha differensiallab quyidagi integral tenglamaga kelamiz:

$$\tau_2'(\eta) \left\{ \lim_{\eta \rightarrow y} \left(P(y, \eta) + \sum_{n=-\infty}^{+\infty} \frac{1}{\sqrt{y-\eta}} e^{-\frac{(2n)}{\sqrt{y-\eta}}} \right) \right\} +$$

$$+ \int_0^y \tau_2'(\eta) \frac{\partial}{\partial y} \left(P(y, \eta) + \sum_{n=-\infty}^{+\infty} \frac{1}{\sqrt{y-\eta}} e^{-\frac{(2n)}{\sqrt{y-\eta}}} \right) d\eta = \Phi_2(y).$$

$$\lim_{\eta \rightarrow y} \left(P(y, \eta) + \sum_{n=-\infty}^{+\infty} \frac{1}{\sqrt{y-\eta}} e^{-\frac{(2n)}{\sqrt{y-\eta}}} \right) = \Phi(y) \text{ belgilash kiritib, } \Phi(y) \neq 0 \text{ deb olib}$$

tenglamani har ikki tomoniga bo'lib yuborsak

$$\tau_2'(\eta) + \int_0^y \tau_2'(\eta) \frac{\frac{\partial}{\partial y} \left(P(y, \eta) + \sum_{n=-\infty}^{+\infty} \frac{1}{\sqrt{y-\eta}} e^{-\frac{(2n)}{\sqrt{y-\eta}}} \right)}{\Phi(y)} d\eta = \frac{\Phi_2(y)}{\Phi(y)}$$

hosil bo'ladi.

$$\frac{\frac{\partial}{\partial y} \left(P(y, \eta) + \sum_{n=-\infty}^{+\infty} \frac{1}{\sqrt{y-\eta}} e^{-\frac{(2n)}{\sqrt{y-\eta}}} \right)}{\Phi(y)} = Q(y, \eta); \quad \frac{\Phi_2(y)}{\Phi(y)} = \Phi(y);$$

belgilashlarni kiritib olamiz.

$Q(y, \eta)$ - yadro uzluksiz yoki 1 dan kichik maxsuslikka ega, $\Phi(y)$ -funksiya esa uzluksiz bo'lsa, $\tau_2'(\eta) + \int_0^y \tau_2'(\eta) Q(y, \eta) d\eta = \Phi(y)$ Volterra 2-tur integral tenglama quyidagi yechimga ega:

$$\tau_2'(\eta) = \Phi(y) + \int_0^y \Phi(y) R_2(y, \eta) d\eta.$$

$R_2(y, \eta)$ - $Q(y, \eta)$ yadroning rezolventasi.

$\tau_2'(y)$ ni (8) tenglikka olib borib qo'yib $v_2^-(x)$ ni topish mumkin bo'ladi.

Topilganlarga asosan qo'yilgan $\{(1),(2),(3),(4),(5)\}$ masala yechimini Ω_0 sohada 1-chegaraviy masala yechimi ko'rinishda Ω_1 va Ω_2 da esa Dalamber formulasi orqali tiklab ko'rsatish mumkin.

Shunday qilib quyidagi tasdiq o'rinli ekanligini isbotladik:

Teorema. Agar

$$K(x, z) = K_1(x)K_1(z), \frac{\partial}{\partial z} K(x, z) = -K_1(x)K_1(z), K(x, x) \geq 0, \int_0^1 \int_0^\xi R(1, z) dz d\xi \neq -1$$

$$, \frac{\partial}{\partial z} P(y, z) = -P_1(y)P_1(z), P(y, y) \geq 0,$$

$$\lim_{\eta \rightarrow y} \left(P(y, \eta) + \sum_{n=-\infty}^{+\infty} \frac{1}{\sqrt{y-\eta}} e^{-\frac{(2n)}{\sqrt{y-\eta}}} \right) = P(y) \neq 0, \varphi(y), \psi_i(x) \in C^1[0, 1], (i=1, 2),$$

$$I_1(f(x)) = \int_0^x f(z)K(x, z)dz, I_2(g(y)) = \int_0^y g(y)P(y, z)dz$$

shartlar bajarilsa, u holda masalaning yechimi mavjud va yagona.

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